# Differential Semblance Migration Velocity Analysis via Reverse Time Migration: Gradient Computation 

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## Outline

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## Introduction

- Differential Semblance Velocity Analysis
- Smoothness
- Convexity
- Reverse Time Migration
- Dip limitation
- DSMVA-RTM

Acoustic wave equation with constant density

$$
\left(\frac{1}{c^{2}(\mathbf{x})} \frac{\partial^{2}}{\partial t^{2}}-\nabla^{2}\right) p\left(\mathbf{x}, t ; \mathbf{x}_{s}\right)=f(t) \delta\left(\mathbf{x}-\mathbf{x}_{s}\right)
$$

$\mathbf{x}=$ position vector
$\mathbf{x}_{\mathbf{s}}=$ position of the point source
$c(\mathbf{x})=$ velocity
$p\left(\mathbf{x}, t ; \mathbf{x}_{\mathbf{s}}\right)=$ pressure
Source time function $f(t)=\delta(t)$
Two-way wave equation operator $\mathbf{L}:=\frac{1}{c^{2}(\mathbf{x})} \frac{\partial^{2}}{\partial t^{2}}-\nabla^{2}$

Source Wavefields (forward in time):

$$
\begin{aligned}
\mathbf{L} S\left(\mathbf{x}, t ; \mathbf{x}_{s}\right) & =\delta(t) \delta\left(\mathbf{x}-\mathbf{x}_{s}\right) \\
S & \equiv 0 \text { for } t<0
\end{aligned}
$$

Receiver Wavefields (backwards in time):

$$
\begin{aligned}
\mathbf{L} R\left(\mathbf{x}, t ; \mathbf{x}_{s}\right) & =\int d\left(\mathbf{x}_{r}, t ; \mathbf{x}_{s}\right) \delta\left(\mathbf{x}-\mathbf{x}_{r}\right) d \mathbf{x}_{r} \\
R & \equiv 0 \text { for } t>t_{\max }
\end{aligned}
$$

## Imaging

Write

$$
c(\mathbf{x})=v(\mathbf{x})+\delta v(\mathbf{x}) .
$$

Forward Born Modeling:

$$
F[v] \delta v=\left.\delta S\right|_{\text {surface }}
$$

Adjoint of Forward Modeling:

$$
F^{*} d=I
$$

Migration formula:

$$
I(\mathbf{x}, h)=\int S\left(\mathbf{x}+h, t ; \mathbf{x}_{s}\right) R\left(\mathbf{x}-h, t ; \mathbf{x}_{s}\right) d \mathbf{x}_{s} d t
$$

DSMVA-RTM objective function:

$$
J[v]=\frac{1}{2}\|P I\|^{2}
$$

## Gradient Computation

For $u=P^{*} P I, D I^{*} u$ gives the gradient.
Introducing $g_{r}$ solving

$$
\begin{aligned}
\mathbf{L} g_{r}\left(\mathbf{x}, t ; \mathbf{x}_{s}\right) & =\int R\left(\mathbf{x}-2 h, t ; \mathbf{x}_{s}\right) u(t-h, h) d h \\
g_{r} & \equiv 0 \text { for } t>t_{\max }, B C
\end{aligned}
$$

Introducing $g_{s}$ solving

$$
\begin{aligned}
\mathbf{L} g_{s}\left(\mathbf{x}, t ; \mathbf{x}_{s}\right) & =\int S\left(\mathbf{x}+2 h, t ; \mathbf{x}_{s}\right) u(t+h, h) d h \\
g_{s} & \equiv 0 \text { for } t<0, B C
\end{aligned}
$$

The gradient of DSMVA-RTM is

$$
\nabla_{v} J=\int \frac{2}{v^{3}}\left(\frac{\partial^{2} S}{\partial t^{2}} g_{r}+\frac{\partial^{2} R}{\partial t^{2}} g_{s}\right) d \mathbf{x}_{s} d t
$$

## Inversion procedure



## Comparison between one-way and two-way WEMVA

One-way wave equation migration:

- Pros: low computational cost
- Cons: dip limitation

Two-way wave equation migration:

- Pros: no dip limitation
- Cons: high computational cost
- Cons: amplitude correction required
- Amplitude Correction (see Rami's work)
- Modification of existing RTM code based on SEAMX, TSOpt to compute objective function and to verify its convexity
- Future modification to create gradient
- Comparison with DS based on downward continuous extrapolation

Perturbation of objective function:

$$
\begin{aligned}
\delta J & =\frac{1}{2}\langle\delta(P I), P I\rangle+\frac{1}{2}\langle P I, \delta(P I)\rangle \\
& \Longrightarrow \delta J=\left\langle\delta c, \operatorname{Re}\left((D I)^{*} P^{*} P I\right)\right\rangle
\end{aligned}
$$

Gradient of objective function:

$$
\nabla_{c} J=\operatorname{Re}\left\{(D I)^{*}\left(P^{*} P I\right)\right\}
$$

For arbitrary $u(x, h)$, we have

$$
\langle\delta I, u\rangle=\langle D I \delta c, u\rangle=\left\langle\delta c, D I^{*} u\right\rangle
$$

When $u=P^{*} P I, D I^{*} u$ gives the gradient.

$$
\begin{aligned}
\langle\delta I, u\rangle & =\int \overline{\delta I} u(x, h) d x d h \\
& =\int \delta S\left(x+h, x_{s}, t\right) \bar{R}\left(x-h, x_{s}, t\right) u(x, h) d x d x_{s} d t d h \\
& +\int S\left(x+h, x_{s}, t\right) \overline{\delta R}\left(x-h, x_{s}, t\right) u(x, h) d x d x_{s} d t d h \\
& =\int \delta S\left(x, x_{s}, t\right)\left\{\int \bar{R}\left(x-2 h, x_{s}, t\right) u(x-h, h) d h\right\} d x d x_{s} d t \\
& +\int \overline{\delta R}\left(x, x_{s}, t\right)\left\{\int S\left(x+2 h, x_{s}, t\right) u(x+h, h) d h\right\} d x d x_{s} d t
\end{aligned}
$$

Introducing $g_{r}$ solving

$$
\begin{aligned}
\mathbf{L} g_{r}\left(\mathbf{x}, t ; \mathbf{x}_{s}\right) & =\int R\left(\mathbf{x}-2 h, t ; \mathbf{x}_{s}\right) u(t-h, h) d h \\
\int_{\Sigma} \nabla S g_{r}-S \nabla g_{r} & =0 \\
\frac{\partial S}{\partial t} g_{r}-S \frac{\partial g_{r}}{\partial t}\left(\mathbf{x}, 0 ; \mathbf{x}_{s}\right) & =0 \\
\frac{\partial S}{\partial t} g_{r}-S \frac{\partial g_{r}}{\partial t}\left(\mathbf{x}, t_{\max } ; \mathbf{x}_{s}\right) & =0
\end{aligned}
$$

Introducing $g_{s}$ solving

$$
\begin{aligned}
\mathbf{L} g_{s}\left(\mathbf{x}, t ; \mathbf{x}_{s}\right) & =\int S\left(\mathbf{x}+2 h, t ; \mathbf{x}_{s}\right) u(t+h, h) d h \\
\int_{\Sigma} \nabla R g_{s}-R \nabla g_{s} & =0 \\
\frac{\partial R}{\partial t} g_{s}-R \frac{\partial g_{s}}{\partial t}\left(\mathbf{x}, 0 ; \mathbf{x}_{s}\right) & =0 \\
\frac{\partial R}{\partial t} g_{s}-R \frac{\partial g_{s}}{\partial t}\left(\mathbf{x}, t_{\max } ; \mathbf{x}_{s}\right) & =0
\end{aligned}
$$

Then,

$$
\begin{aligned}
\langle\delta I, u\rangle & =\int\left(\delta S \mathbf{L} g_{r}+\delta R \mathbf{L} g_{s}\right) d x d x_{s} d t \\
& =\int\left((\mathbf{L} \delta S) g_{r}+(\mathbf{L} \delta R) g_{s}\right) d x d x_{s} d t \\
& =\int \frac{2 \delta c}{c^{3}} \frac{\partial^{2} S}{\partial t^{2}} g_{r}+\frac{2 \delta c}{c^{3}} \frac{\partial^{2} R}{\partial t^{2}} g_{s} d x d x_{s} d t
\end{aligned}
$$

Thus

$$
\nabla_{c} J=\int \frac{2}{c^{3}}\left(\frac{\partial^{2} S}{\partial t^{2}} g_{r}+\frac{\partial^{2} R}{\partial t^{2}} g_{s}\right) d x_{s} d t
$$

