

Multivalued traveltimes via Liouville equations

Recent developments in level set methods

Jianliang Qian

TRIP, Rice University

Department of Mathematics, UCLA

TRIP Annual Meeting, Houston, Texas

Friday 29 October 2004

Outline

- Overview for multivalued geometrical optics

Outline

- Overview for multivalued geometrical optics
- Liouville equations for geometrical optics

Outline

- Overview for multivalued geometrical optics
- Liouville equations for geometrical optics
- Pseudo-spectral discontinuous Galerkin formulation

Outline

- Overview for multivalued geometrical optics
- Liouville equations for geometrical optics
- Pseudo-spectral discontinuous Galerkin formulation
- Paraxial formulation

Outline

- Overview for multivalued geometrical optics
- Liouville equations for geometrical optics
- Pseudo-spectral discontinuous Galerkin formulation
- Paraxial formulation
- What's next

Eikonal eqn: Hamilton-Jacobi Eqns

- High frequency asymptotics for wave equations:

$$H(\mathbf{x}, \mathbf{p}) = H(\mathbf{x}, \nabla\tau) = 1$$

Eikonal eqn: Hamilton-Jacobi Eqns

- High frequency asymptotics for wave equations:

$$H(\mathbf{x}, \mathbf{p}) = H(\mathbf{x}, \nabla\tau) = 1$$

- Isotropic eikonal equation:

$$c(\mathbf{x})|\nabla\tau| = 1$$

Eikonal eqn: Hamilton-Jacobi Eqns

- High frequency asymptotics for wave equations:

$$H(\mathbf{x}, \mathbf{p}) = H(\mathbf{x}, \nabla\tau) = 1$$

- Isotropic eikonal equation:

$$c(\mathbf{x})|\nabla\tau| = 1$$

- H is a homogeneous Hamiltonian of degree one in p .

Overview: Eulerian methods

- Ray tracing (Lagrangian) based methods:

Overview: Eulerian methods

- Ray tracing (Lagrangian) based methods:
 - Advantage: multi-valued solutions.

Overview: Eulerian methods

- Ray tracing (Lagrangian) based methods:
 - Advantage: multi-valued solutions.
 - Disadvantage: nonuniform coverage of domains.

Overview: Eulerian methods

- Ray tracing (Lagrangian) based methods:
 - Advantage: multi-valued solutions.
 - Disadvantage: nonuniform coverage of domains.
- Finite-difference (Eulerian) based eikonal solvers:

Overview: Eulerian methods

- Ray tracing (Lagrangian) based methods:
 - Advantage: multi-valued solutions.
 - Disadvantage: nonuniform coverage of domains.
- Finite-difference (Eulerian) based eikonal solvers:
 - Advantage: Uniform coverage of domains.

Overview: Eulerian methods

- Ray tracing (Lagrangian) based methods:
 - Advantage: multi-valued solutions.
 - Disadvantage: nonuniform coverage of domains.
- Finite-difference (Eulerian) based eikonal solvers:
 - Advantage: Uniform coverage of domains.
 - Disadvantage: first arrivals only.

Overview: Eulerian methods

- Ray tracing (Lagrangian) based methods:
 - Advantage: multi-valued solutions.
 - Disadvantage: nonuniform coverage of domains.
- Finite-difference (Eulerian) based eikonal solvers:
 - Advantage: Uniform coverage of domains.
 - Disadvantage: first arrivals only.
- Challenge: Eulerian method for multi-valued travel-times

Overview: Eulerian methods

- Ray tracing (Lagrangian) based methods:
 - Advantage: multi-valued solutions.
 - Disadvantage: nonuniform coverage of domains.
- Finite-difference (Eulerian) based eikonal solvers:
 - Advantage: Uniform coverage of domains.
 - Disadvantage: first arrivals only.
- Challenge: Eulerian method for multi-valued travel-times
- Caustics decomposition (Benamou'99), slowness matching (Symes'96), segment projection (Engquist, etal '02), level sets (Osher, etal'02, Fomel-Sethian'02)

Lagrangian ray-tracing equations



$$H(\mathbf{x}, \mathbf{p}) = H(\mathbf{x}, \nabla\tau) = 1$$

Lagrangian ray-tracing equations



$$H(\mathbf{x}, \mathbf{p}) = H(\mathbf{x}, \nabla\tau) = 1$$

- Hamiltonian system:

$$\frac{d\tilde{\mathbf{x}}}{dt} = \nabla_{\tilde{\mathbf{p}}} H(\tilde{\mathbf{x}}, \tilde{\mathbf{p}}), \quad \tilde{\mathbf{x}}(0, \mathbf{x}, \mathbf{p}) = \mathbf{x}$$

$$\frac{d\tilde{\mathbf{p}}}{dt} = -\nabla_{\tilde{\mathbf{x}}} H(\tilde{\mathbf{x}}, \tilde{\mathbf{p}}), \quad \tilde{\mathbf{p}}(0, \mathbf{x}, \mathbf{p}) = \mathbf{p}$$

$$\frac{d\tilde{\tau}}{dt} = \tilde{\mathbf{p}} \cdot \nabla_{\tilde{\mathbf{p}}} H(\tilde{\mathbf{x}}, \tilde{\mathbf{p}}), \quad \tilde{\tau}(0, \mathbf{x}, \mathbf{p}) = 0.$$

Liouville equations

$$\begin{aligned}\frac{d\tilde{\mathbf{x}}}{dt} &= \nabla_{\mathbf{p}}H(\tilde{\mathbf{x}}, \tilde{\mathbf{p}}), & \tilde{\mathbf{x}}(0, \mathbf{x}, \mathbf{p}) &= \mathbf{x} \\ \frac{d\tilde{\mathbf{p}}}{dt} &= -\nabla_{\mathbf{x}}H(\tilde{\mathbf{x}}, \tilde{\mathbf{p}}), & \tilde{\mathbf{p}}(0, \mathbf{x}, \mathbf{p}) &= \mathbf{p}\end{aligned}$$

Liouville equations



$$\begin{aligned}\frac{d\tilde{\mathbf{x}}}{dt} &= \nabla_{\mathbf{p}}H(\tilde{\mathbf{x}}, \tilde{\mathbf{p}}), & \tilde{\mathbf{x}}(0, \mathbf{x}, \mathbf{p}) &= \mathbf{x} \\ \frac{d\tilde{\mathbf{p}}}{dt} &= -\nabla_{\mathbf{x}}H(\tilde{\mathbf{x}}, \tilde{\mathbf{p}}), & \tilde{\mathbf{p}}(0, \mathbf{x}, \mathbf{p}) &= \mathbf{p}\end{aligned}$$

- defines a Hamiltonian flow \mathcal{F}_t

$$\mathcal{F}_t(\mathbf{x}, \mathbf{p}) = (\tilde{\mathbf{x}}(t, \mathbf{x}, \mathbf{p}), \tilde{\mathbf{p}}(t, \mathbf{x}, \mathbf{p}))$$

Liouville equations



$$\begin{aligned}\frac{d\tilde{\mathbf{x}}}{dt} &= \nabla_{\mathbf{p}}H(\tilde{\mathbf{x}}, \tilde{\mathbf{p}}), & \tilde{\mathbf{x}}(0, \mathbf{x}, \mathbf{p}) &= \mathbf{x} \\ \frac{d\tilde{\mathbf{p}}}{dt} &= -\nabla_{\mathbf{x}}H(\tilde{\mathbf{x}}, \tilde{\mathbf{p}}), & \tilde{\mathbf{p}}(0, \mathbf{x}, \mathbf{p}) &= \mathbf{p}\end{aligned}$$

- defines a Hamiltonian flow \mathcal{F}_t

$$\mathcal{F}_t(\mathbf{x}, \mathbf{p}) = (\tilde{\mathbf{x}}(t, \mathbf{x}, \mathbf{p}), \tilde{\mathbf{p}}(t, \mathbf{x}, \mathbf{p}))$$

- associated to the Liouville equation

$$w_t + \nabla_{\mathbf{p}}H \cdot \nabla_{\mathbf{x}}w - \nabla_{\mathbf{x}}H \cdot \nabla_{\mathbf{p}}w = 0$$

Liouville equations

$$\begin{aligned}\frac{d\tilde{\mathbf{x}}}{dt} &= \nabla_{\mathbf{p}}H(\tilde{\mathbf{x}}, \tilde{\mathbf{p}}), & \tilde{\mathbf{x}}(0, \mathbf{x}, \mathbf{p}) &= \mathbf{x} \\ \frac{d\tilde{\mathbf{p}}}{dt} &= -\nabla_{\mathbf{x}}H(\tilde{\mathbf{x}}, \tilde{\mathbf{p}}), & \tilde{\mathbf{p}}(0, \mathbf{x}, \mathbf{p}) &= \mathbf{p}\end{aligned}$$

- defines a Hamiltonian flow \mathcal{F}_t

$$\mathcal{F}_t(\mathbf{x}, \mathbf{p}) = (\tilde{\mathbf{x}}(t, \mathbf{x}, \mathbf{p}), \tilde{\mathbf{p}}(t, \mathbf{x}, \mathbf{p}))$$

- associated to the Liouville equation

$$w_t + \nabla_{\mathbf{p}}H \cdot \nabla_{\mathbf{x}}w - \nabla_{\mathbf{x}}H \cdot \nabla_{\mathbf{p}}w = 0$$

- Liouville equation shares the same bicharacteristics as the original nonlinear PDE.

Properties of Liouville equations

$$w_t + \nabla_{\mathbf{p}}H \cdot \nabla_{\mathbf{x}}w - \nabla_{\mathbf{x}}H \cdot \nabla_{\mathbf{p}}w = 0$$

- A linear transport equation for w in phase space.

Properties of Liouville equations

$$w_t + \nabla_{\mathbf{p}}H \cdot \nabla_{\mathbf{x}}w - \nabla_{\mathbf{x}}H \cdot \nabla_{\mathbf{p}}w = 0$$

- A linear transport equation for w in phase space.
- If the Hamiltonian is smooth, then a global solution exists and is unique for all t .

Properties of Liouville equations

$$w_t + \nabla_{\mathbf{p}}H \cdot \nabla_{\mathbf{x}}w - \nabla_{\mathbf{x}}H \cdot \nabla_{\mathbf{p}}w = 0$$

- A linear transport equation for w in phase space.
- If the Hamiltonian is smooth, then a global solution exists and is unique for all t .
- Liouville: solution remains constant along all bicharacteristics $(\tilde{\mathbf{x}}, \tilde{\mathbf{p}})$.

Properties of Liouville equations

$$w_t + \nabla_{\mathbf{p}}H \cdot \nabla_{\mathbf{x}}w - \nabla_{\mathbf{x}}H \cdot \nabla_{\mathbf{p}}w = 0$$

- A linear transport equation for w in phase space.
- If the Hamiltonian is smooth, then a global solution exists and is unique for all t .
- Liouville: solution remains constant along all bicharacteristics $(\tilde{\mathbf{x}}, \tilde{\mathbf{p}})$.
- No caustics appears in phase space; singularities are unfolded in phase space.

Properties of Liouville equations

$$w_t + \nabla_{\mathbf{p}}H \cdot \nabla_{\mathbf{x}}w - \nabla_{\mathbf{x}}H \cdot \nabla_{\mathbf{p}}w = 0$$

- A linear transport equation for w in phase space.
- If the Hamiltonian is smooth, then a global solution exists and is unique for all t .
- Liouville: solution remains constant along all bicharacteristics $(\tilde{\mathbf{x}}, \tilde{\mathbf{p}})$.
- No caustics appears in phase space; singularities are unfolded in phase space.
- Nonlinear first order eqn: local smooth solutions only.

Liouville eqns: isotropic media

- Isotropic eikonal equation:

$$c|\nabla\tau| = 1$$

$$H(\mathbf{x}, \mathbf{p}) = c(\mathbf{x})|\mathbf{p}| = 1$$

Liouville eqns: isotropic media

- Isotropic eikonal equation:

$$c|\nabla\tau| = 1$$

$$H(\mathbf{x}, \mathbf{p}) = c(\mathbf{x})|\mathbf{p}| = 1$$

- Liouville equation:

$$w_t + c(\mathbf{x})\frac{\mathbf{p}}{|\mathbf{p}|} \cdot \nabla_{\mathbf{x}}w - \nabla_{\mathbf{x}}(c(\mathbf{x})|\mathbf{p}|) \cdot \nabla_{\mathbf{p}}w = 0$$

Liouville eqns: isotropic media

- Isotropic eikonal equation:

$$c|\nabla\tau| = 1$$

$$H(\mathbf{x}, \mathbf{p}) = c(\mathbf{x})|\mathbf{p}| = 1$$

- Liouville equation:

$$w_t + c(\mathbf{x})\frac{\mathbf{p}}{|\mathbf{p}|} \cdot \nabla_{\mathbf{x}}w - \nabla_{\mathbf{x}}(c(\mathbf{x})|\mathbf{p}|) \cdot \nabla_{\mathbf{p}}w = 0$$

- Reduced form (Engquist-Runborg'02):

$$w(t, \mathbf{x}, \mathbf{p}) = c(\mathbf{x})\delta\left(|\mathbf{p}| - \frac{1}{c(\mathbf{x})}\right)\mathbf{u}\left(\mathbf{t}, \mathbf{x}, \frac{\mathbf{p}}{|\mathbf{p}|}\right)$$

Isotropic Liouville: cont.

- Reduced form (Engquist-Runborg'02):

$$w(t, \mathbf{x}, \mathbf{p}) = c(\mathbf{x}) \delta\left(|\mathbf{p}| - \frac{1}{c(\mathbf{x})}\right) \mathbf{u}\left(t, \mathbf{x}, \frac{\mathbf{p}}{|\mathbf{p}|}\right)$$

Isotropic Liouville: cont.

- Reduced form (Engquist-Runborg'02):

$$w(t, \mathbf{x}, \mathbf{p}) = c(\mathbf{x}) \delta\left(|\mathbf{p}| - \frac{1}{c(\mathbf{x})}\right) u\left(t, \mathbf{x}, \frac{\mathbf{p}}{|\mathbf{p}|}\right)$$

- Two dimension: $\mathbf{p} = (r \cos \theta, r \sin \theta)$.

$$u_t + c \cos(\theta) u_{x_1} + c \sin(\theta) u_{x_2} \\ + (c_{x_1} \sin(\theta) - c_{x_2} \cos(\theta)) u_\theta = 0$$

in a reduced phase space (x_1, x_2, θ) .

Isotropic Liouville: cont.

- Reduced form (Engquist-Runborg'02):

$$w(t, \mathbf{x}, \mathbf{p}) = c(\mathbf{x}) \delta\left(|\mathbf{p}| - \frac{1}{c(\mathbf{x})}\right) u\left(t, \mathbf{x}, \frac{\mathbf{p}}{|\mathbf{p}|}\right)$$

- Two dimension: $\mathbf{p} = (r \cos \theta, r \sin \theta)$.

$$u_t + c \cos(\theta) u_{x_1} + c \sin(\theta) u_{x_2} \\ + (c_{x_1} \sin(\theta) - c_{x_2} \cos(\theta)) u_\theta = 0$$

in a reduced phase space (x_1, x_2, θ) .

- A similar equation exists in 3-D case.

Co-dimension m object in \mathcal{R}^n

- Wavefront for the ray tracing system in reduced phase space (\mathbf{x}, \mathbf{s}) :

Co-dimension m object in \mathcal{R}^n

- Wavefront for the ray tracing system in reduced phase space (\mathbf{x}, \mathbf{s}) :
 - 2-D physical space: $\mathbf{x} \in \mathcal{R}^2$, $\frac{\mathbf{p}}{|\mathbf{p}|} \in \mathcal{S}^1 \Rightarrow$
Wavefront of $\dim = 1$ and co-dimension = 2; a curve intersected by 2 level sets

Co-dimension m object in \mathcal{R}^n

- Wavefront for the ray tracing system in reduced phase space (\mathbf{x}, \mathbf{s}) :
 - 2-D physical space: $\mathbf{x} \in \mathcal{R}^2$, $\frac{\mathbf{p}}{|\mathbf{p}|} \in \mathcal{S}^1 \Rightarrow$
Wavefront of $\dim = 1$ and co-dimension = 2; a curve intersected by 2 level sets
 - $\mathbf{x} \in \mathcal{R}^3$, $\frac{\mathbf{p}}{|\mathbf{p}|} \in \mathcal{S}^2 \Rightarrow$ Wavefront of $\dim = 2$ and co-dimension = 3; a surface intersected by 3 level sets

Level set motion equation

- Define scalar level set functions: $\phi^k = \phi^k(t, \mathbf{x}, \mathbf{p})$,
 $k = 1, 2, \dots, m$ (with $m=2$ -D or 3 -D).

Level set motion equation

- Define scalar level set functions: $\phi^k = \phi^k(t, \mathbf{x}, \mathbf{p})$, $k = 1, 2, \dots, m$ (with $m=2$ -D or 3 -D).
- The level set motion equation:

$$\phi_t^k + \nabla_{\mathbf{p}} H \cdot \phi_{\mathbf{x}}^k - \nabla_{\mathbf{x}} H \cdot \phi_{\mathbf{p}}^k = 0, \quad k = 1, 2, \dots, m$$

Level set motion equation

- Define scalar level set functions: $\phi^k = \phi^k(t, \mathbf{x}, \mathbf{p})$, $k = 1, 2, \dots, m$ (with $m=2$ -D or 3 -D).
- The level set motion equation:

$$\phi_t^k + \nabla_{\mathbf{p}} H \cdot \phi_{\mathbf{x}}^k - \nabla_{\mathbf{x}} H \cdot \phi_{\mathbf{p}}^k = 0, \quad k = 1, 2, \dots, m$$

- Compact form (2 or 3 linear eqns):

$$\Phi_t + \nabla_{\mathbf{x}} \Phi \nabla_{\mathbf{p}} H - \nabla_{\mathbf{p}} \Phi \nabla_{\mathbf{x}} H = 0$$

where $\Phi = (\phi^1, \dots, \phi^m)^T$.

Level set motion equation

- Define scalar level set functions: $\phi^k = \phi^k(t, \mathbf{x}, \mathbf{p})$, $k = 1, 2, \dots, m$ (with $m=2$ -D or 3-D).
- The level set motion equation:

$$\phi_t^k + \nabla_{\mathbf{p}} H \cdot \phi_{\mathbf{x}}^k - \nabla_{\mathbf{x}} H \cdot \phi_{\mathbf{p}}^k = 0, \quad k = 1, 2, \dots, m$$

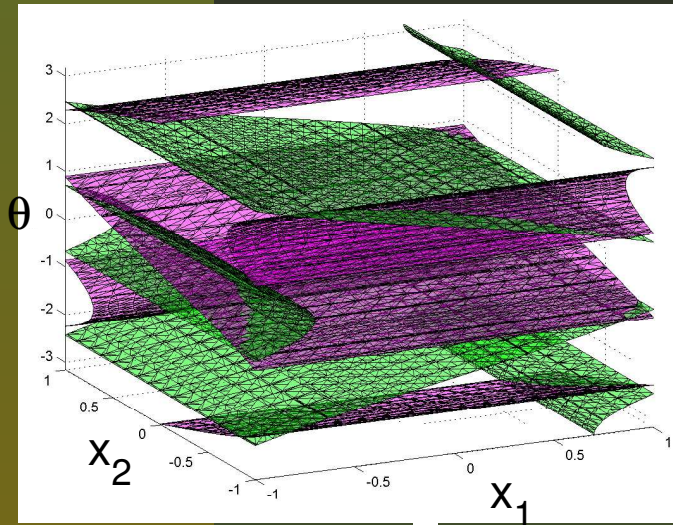
- Compact form (2 or 3 linear eqns):

$$\Phi_t + \nabla_{\mathbf{x}} \Phi \nabla_{\mathbf{p}} H - \nabla_{\mathbf{p}} \Phi \nabla_{\mathbf{x}} H = 0$$

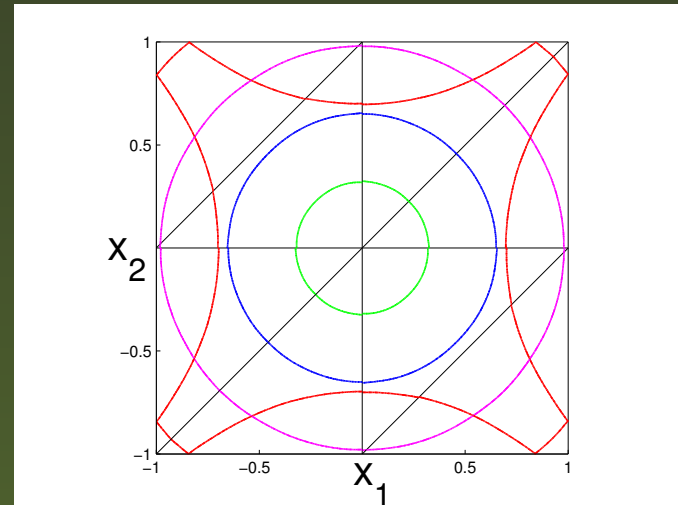
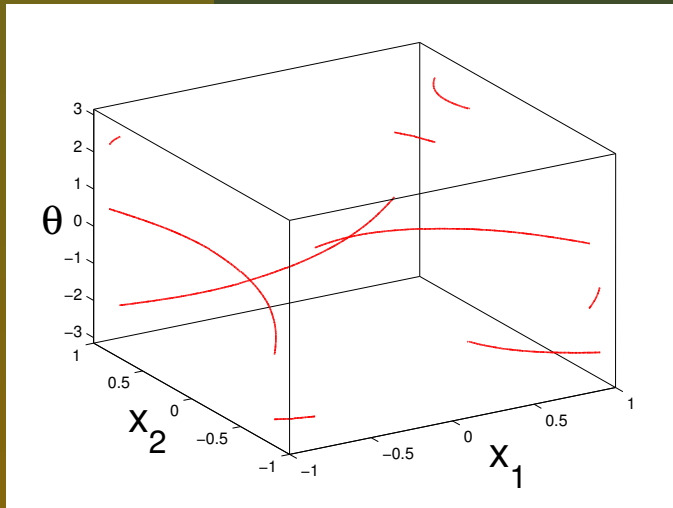
where $\Phi = (\phi^1, \dots, \phi^m)^T$.

- Initialize those equations with initial wavefronts in phase space

Extract multiple traveltimes



- $\Gamma(t) = \{(\mathbf{x}, \mathbf{p}) : \Phi(t, \mathbf{x}, \mathbf{p}) = 0\}$: wavefront in phase space at time t .

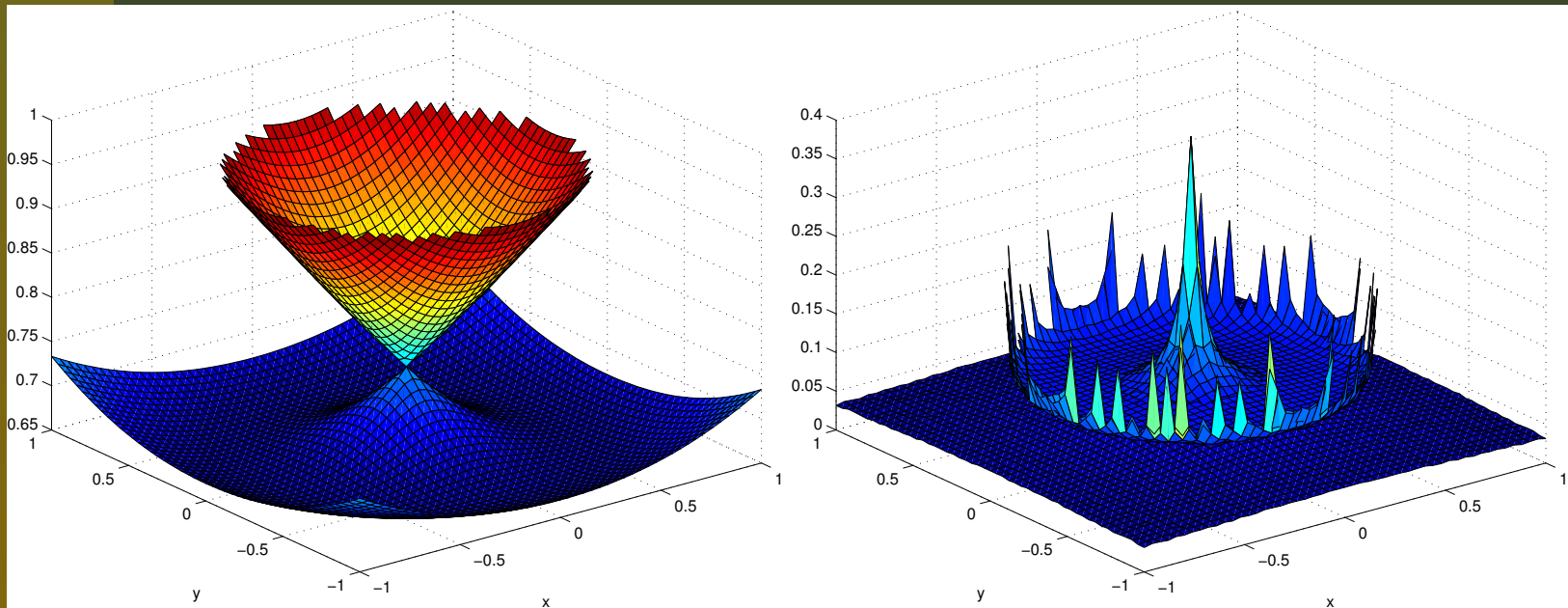


Recent Developments

- Spectral/discontinuous Galerkin (DG) finite-element formulation (Cockburn-Qian-Reitich-Wang'04)

Recent Developments

- Spectral/discontinuous Galerkin (DG) finite-element formulation (Cockburn-Qian-Reitich-Wang'04)
- Paraxial formulation for 3-D geometrical optics (Leung-Qian-Osher'04)



Spectral/DG formulation

- Liouville equation:

$$u_t + c \cos(\theta) u_{x_1} + c \sin(\theta) u_{x_2} + (c_{x_1} \sin(\theta) - c_{x_2} \cos(\theta)) u_\theta = 0$$

- Pseudo-Spectral formulation:

$$u(x_1, x_2, \theta, t) = \sum_{n=-N}^N U_n(x_1, x_2, t) e^{in\theta}$$

Spectral/DG formulation: cont.

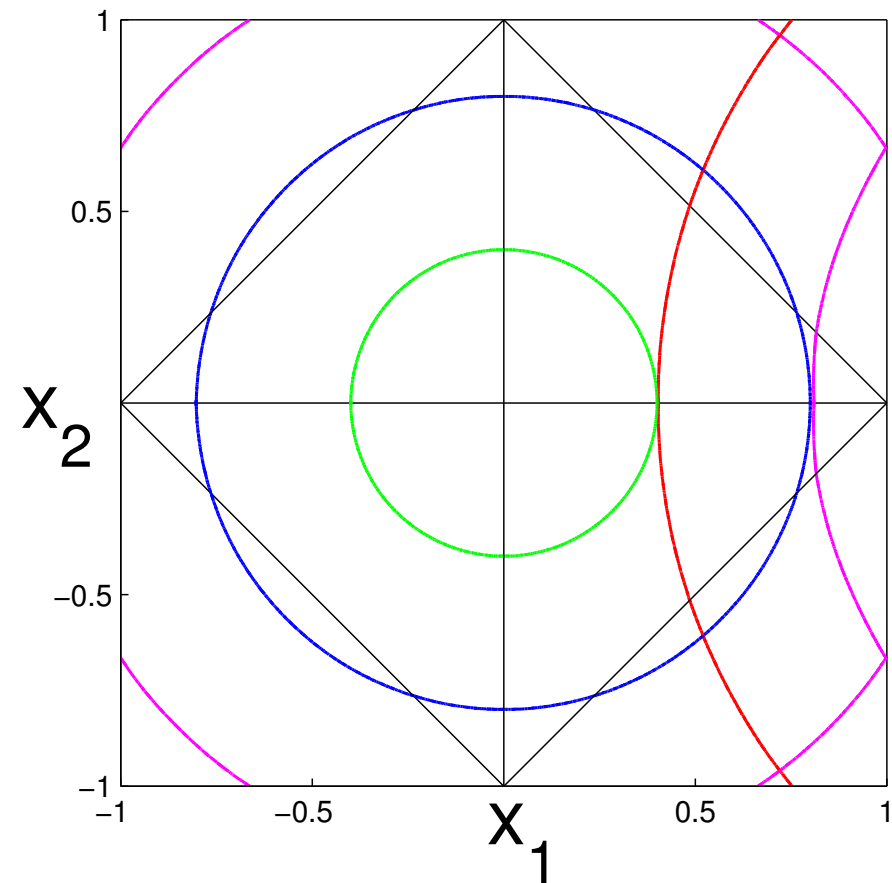
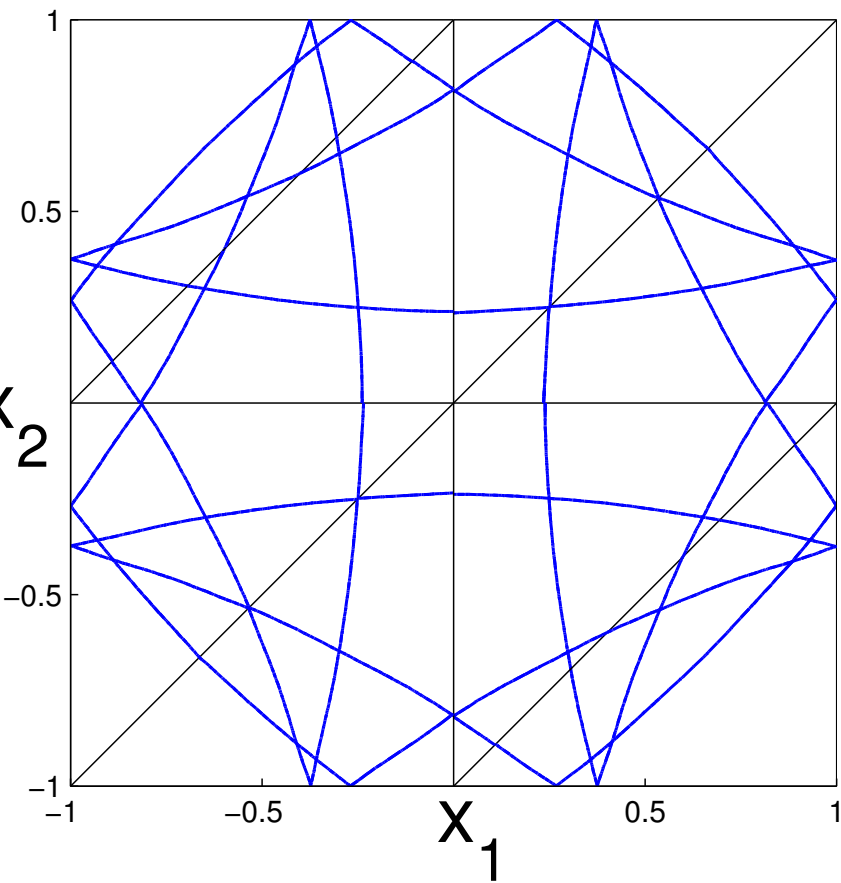
- Strictly symmetric and explicitly diagonalizable hyperbolic system for spectral coefficients:

$$\mathbf{U}_t + A_1 \mathbf{U}_{x_1} + A_2 \mathbf{U}_{x_2} + B\mathbf{U} = 0$$

- Apply discontinuous Galerkin formulation (Cockburn-Shu'01)
- DG: Easily parallelizable and arbitrarily high order accuracy

Spectral/DG formulation: results

- Multiple reflection:



Paraxial Liouville formulation: 3-D

- Paraxial ray tracing system:

$$x_z = \frac{1}{\cos \psi \tan \theta}, \quad y_z = \tan \psi,$$
$$\theta_z = \frac{c_x}{c \cos \psi} - \frac{c_z + c_y \tan \psi}{c \tan \theta},$$
$$\psi_z = \frac{c_z \tan \psi - c_y}{c \sin^2 \theta}$$

Paraxial Liouville formulation: 3-D

- Paraxial ray tracing system:

$$x_z = \frac{1}{\cos \psi \tan \theta}, \quad y_z = \tan \psi,$$
$$\theta_z = \frac{c_x}{c \cos \psi} - \frac{c_z + c_y \tan \psi}{c \tan \theta},$$
$$\psi_z = \frac{c_z \tan \psi - c_y}{c \sin^2 \theta}$$

- Paraxial Liouville eqns for level sets and traveltimes:

$$\phi_z^m + x_z \phi_x^m + y_z \phi_y^m + \theta_z \phi_\theta^m + \psi_z \phi_\psi^m = 0$$

$$T_z + x_z T_x + y_z T_y + \theta_z T_\theta + \psi_z T_\psi = \frac{1}{c \sin \theta \cos \psi}$$

Paraxial Liouville formulation: cont.

- Single source and multiple sources can be treated in the same framework.

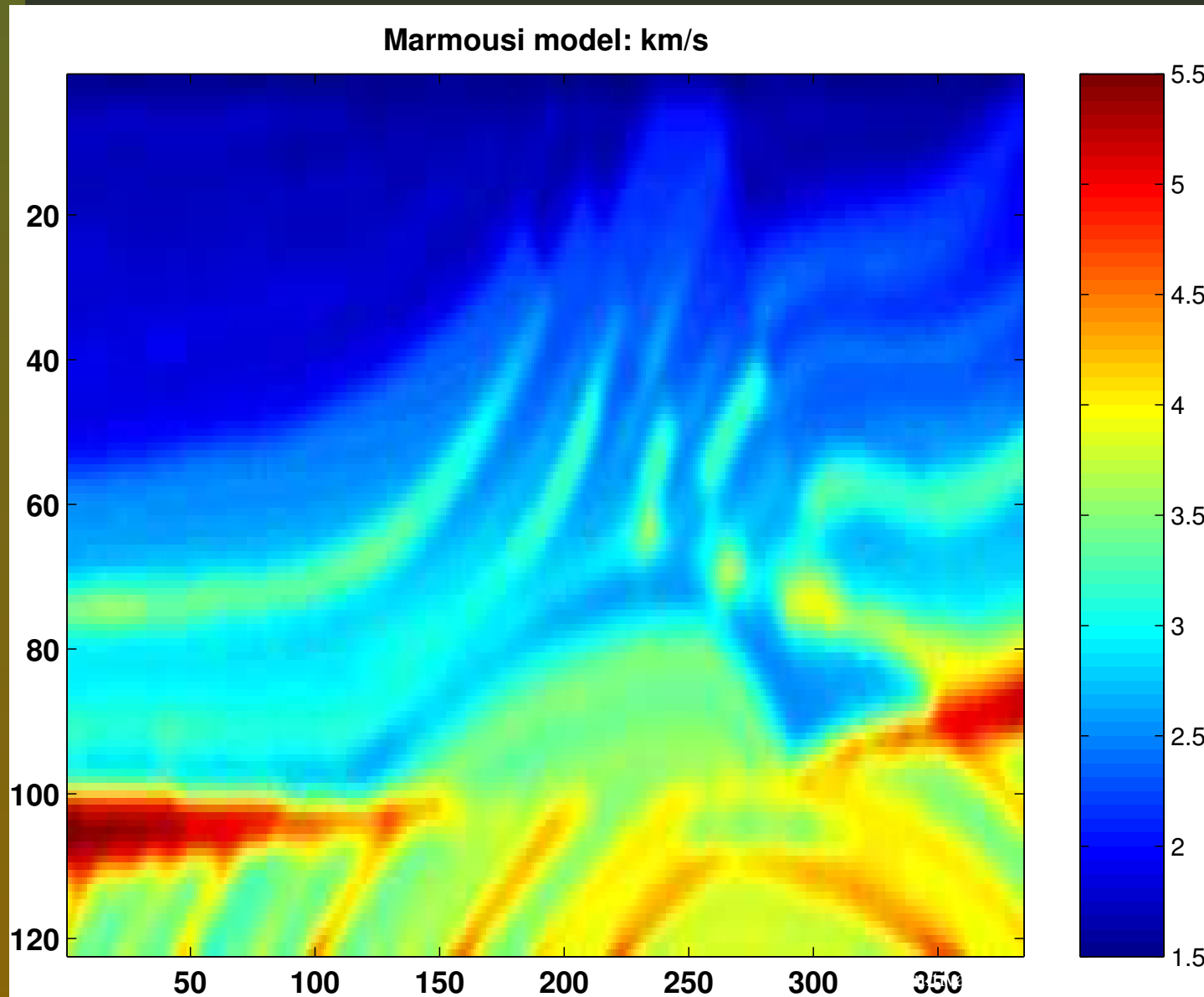
Paraxial Liouville formulation: cont.

- Single source and multiple sources can be treated in the same framework.
- Amplitude can be computed in the same framework

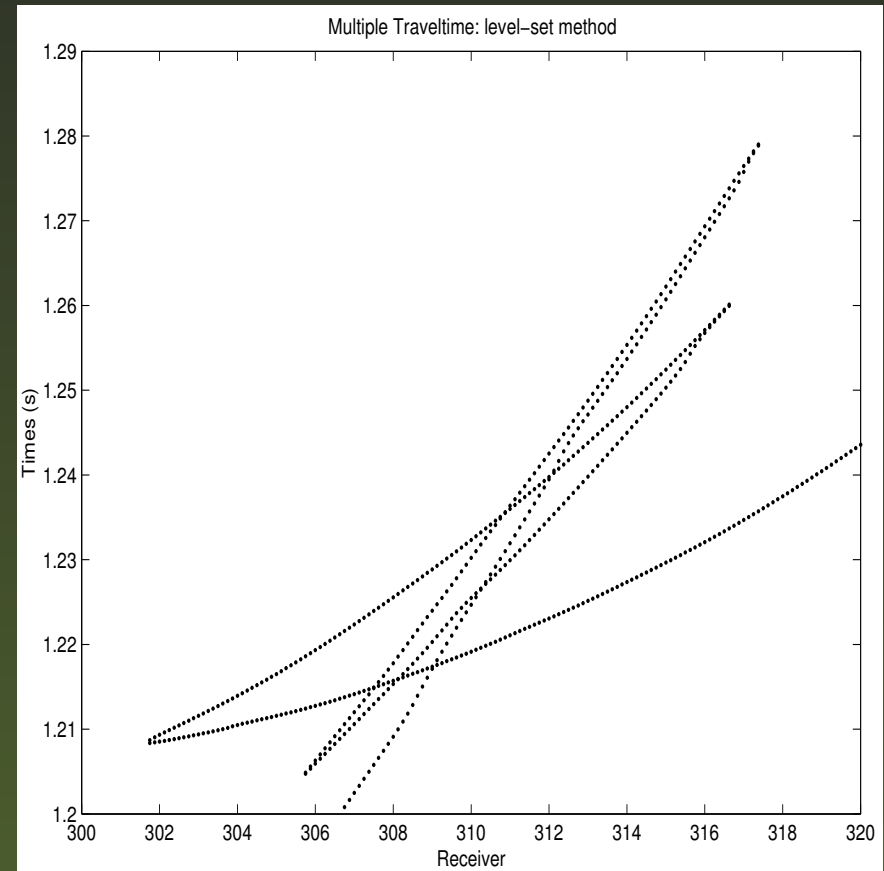
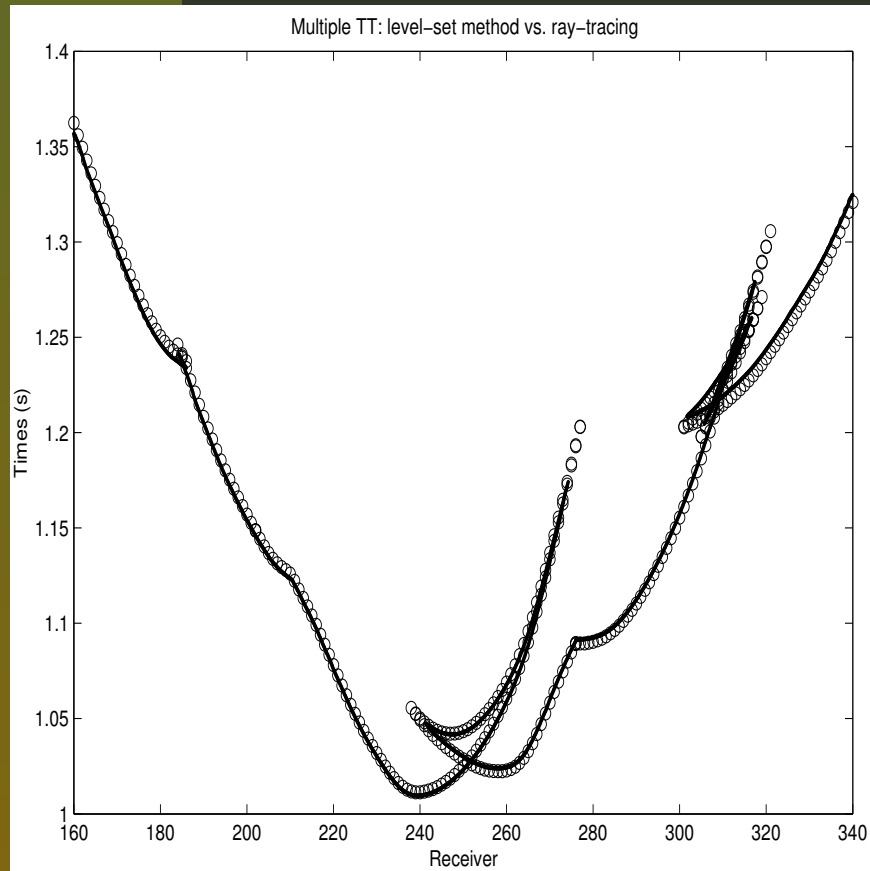
Paraxial Liouville formulation: cont.

- Single source and multiple sources can be treated in the same framework.
- Amplitude can be computed in the same framework
- Efficient implementation by using Semi-Lagrangian method

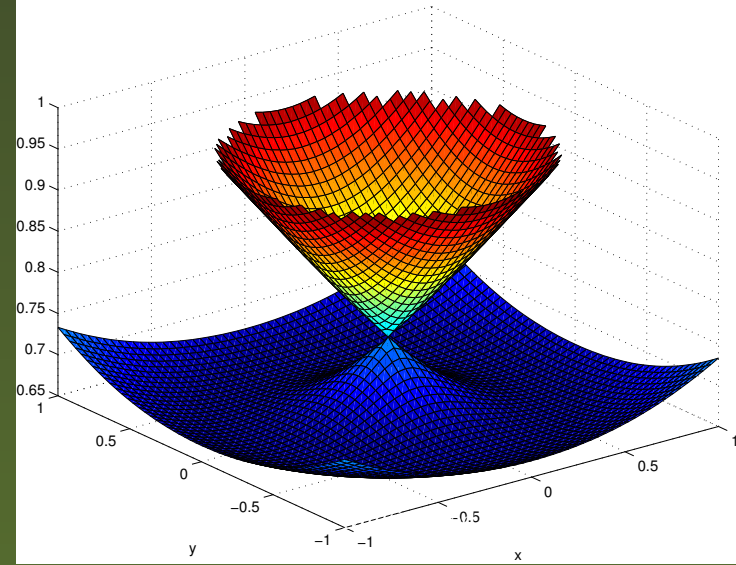
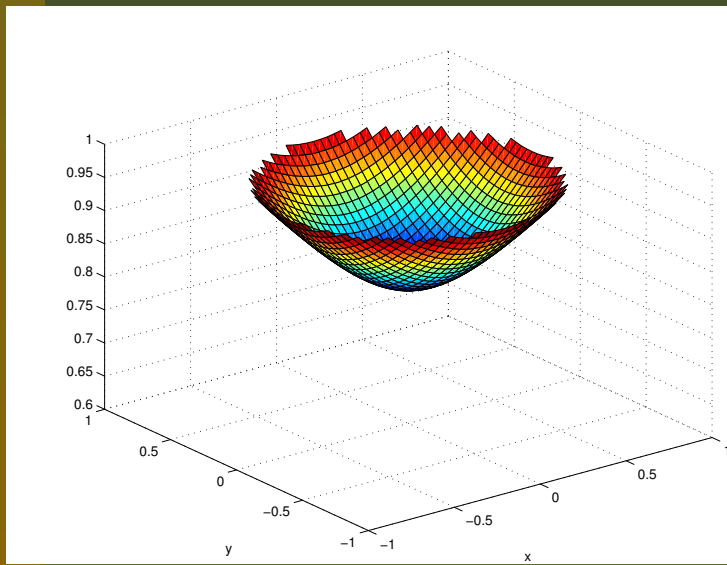
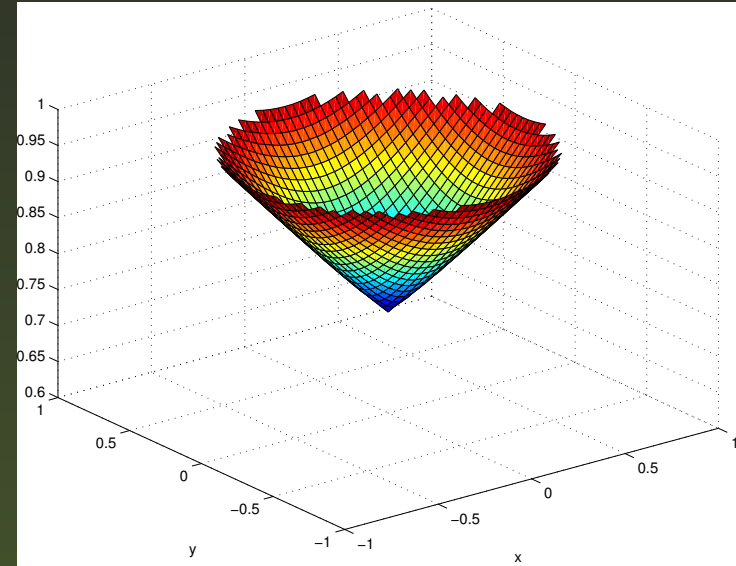
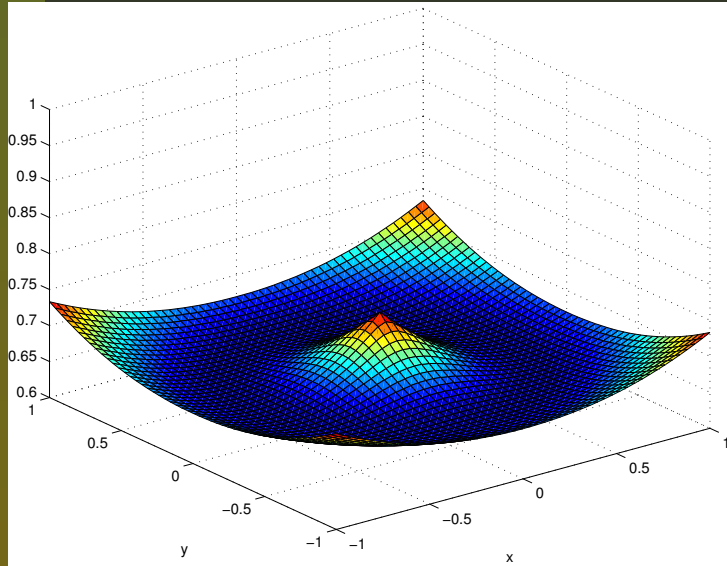
Synthetic Marmousi model



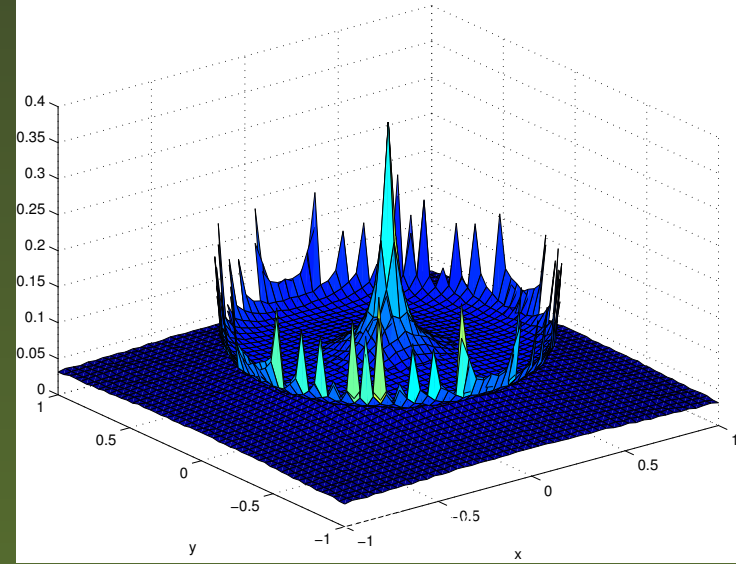
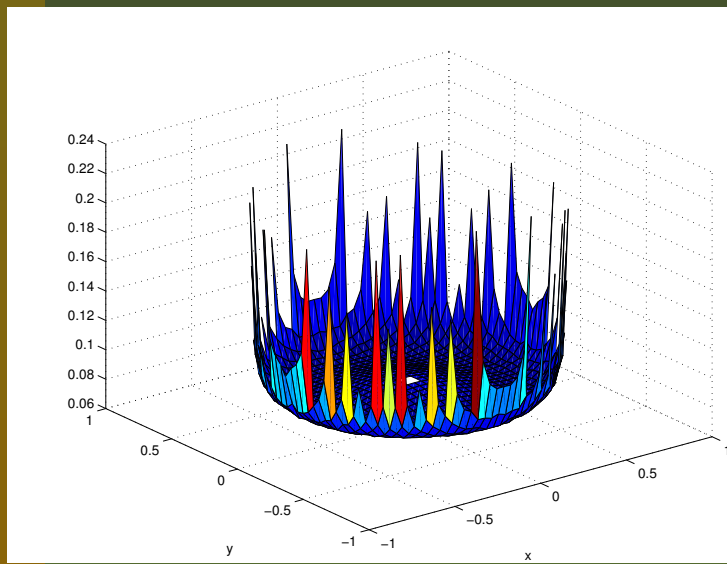
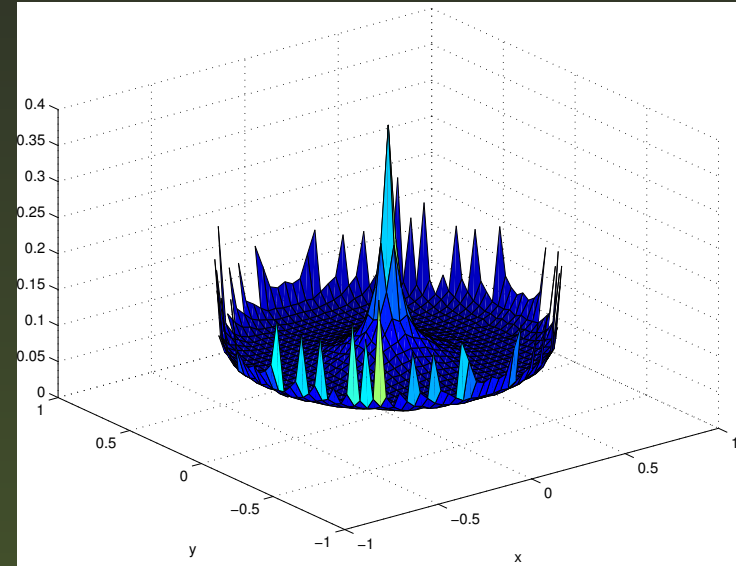
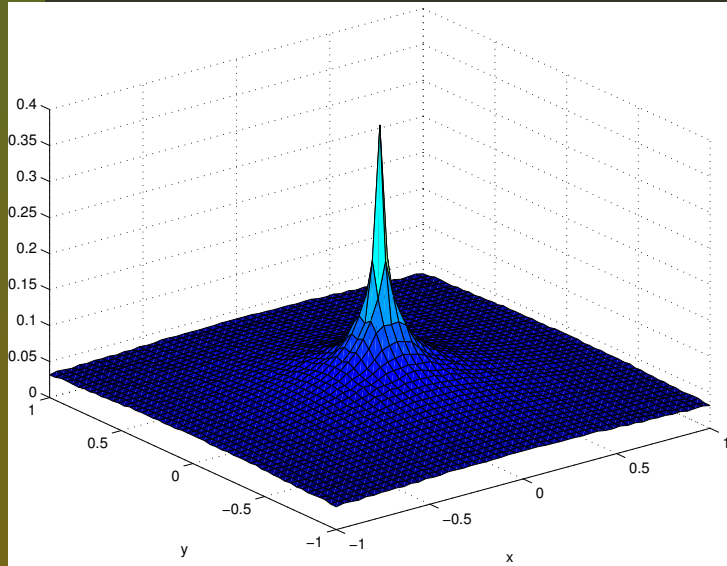
Marmousi: traveltimes



3D Vinje's Gaussian: travelttime



3D Vinje's Gaussian: amplitude



What's next

- Incorporate into seismic migration

What's next

- Incorporate into seismic migration
- Multivalued high resolution reflection tomography